

Optical patterning of magnetic domains and defects in ferromagnetic liquid crystal colloids

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A promising approach in designing composite materials with an unusual physical behavior combines solid nanocrystals and orientationally ordered soft matter materials. Such composites can not only inherit properties of their constituents but also can exhibit emergent behaviors such as ferromagnetic ordering of colloidal metal nanoparticles, forming magnetic monopole domains, when dispersed in a nematic liquid crystal. Here, we demonstrate the optical patterning of domains, crystals, and topological defects in such ferromagnetic liquid crystal colloids, which allow for patterning their response to magnetic fields. Ordering, reorientation of the defects in terms of magnetic moment, differ significantly compared to non-polar nematic and ferromagnetic alike. © 2015 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4928552>]

Liquid crystal (LC) colloids attract considerable attention of scientists in order to benefit from the richness of their

500 mg of 5CB, 95 μl of pentacene (5CB, Chengde
 Yonghua Materials Co. Ltd.) and 5 μl of
 a benzene-concentrated beam 1205 LC (Beam Co.). Non-
 isomerizing FLCs (pencil 2) were used for
 optical defect patterning. 100 μl of pure
 5CB and hexamethylbenzene were used for the
 preparation of the FLC. Silane-PEG capped magnetic
 nanoparticles were added to the LC. 5 μl of
 ethanol was added to 15 μl of the LC mixture to bring
 the isotropic phase, followed by adding 15 μl of 0.1%
 magnetic nanoparticles in ethanol. The sample was kept
 at 90 °C for 3 h to allow the ethanol to evaporate, yielding
 a cell in the isotropic phase. The sample was then
 rapidly cooled to the nematic phase of the mixture
 while rotating at 2000 rpm for 5 min to remove residual
 aggregates. The nematic phase was aligned on
 the substrate by rubbing with a cloth. The
 magnetic nanoparticles were added to the LC mixture
 to bring the isotropic phase, followed by adding 15 μl
 of 0.1% magnetic nanoparticles in ethanol. The sample
 was kept at 90 °C for 3 h to allow the ethanol to
 evaporate, yielding a cell in the isotropic phase. The
 sample was then rapidly cooled to the nematic phase
 of the mixture while rotating at 2000 rpm for 5 min
 to remove residual aggregates. The nematic phase
 was aligned on the substrate by rubbing with a cloth.
 Homeotropic glass cells in the pol domain
 FLCs and M pointing in one of the two in-plane
 directions along the vertical far-field direction
 were prepared using 1- or 0.17-mm thick glass
 plates treated with an aqueous solution of 0.1%
 N,N-dimethyl-N-ocadecyl-3-aminopropylamine
 hydrochloride (DMOAP, Acros Organics) via
 dip-coating. The cell gap thicknesses of 30 μm
 or 60 μm were obtained by spacing the glass
 plates with their edges with UV-curable
 adhesive (NOA-65, Norland Products) or
 epoxy coating. The thicknesses of the
 coating were controlled by controlling the
 diameter. Cell filling was done at room
 temperature using capillary action. FLC
 cells were prepared in a random magnetic
 domain of lateral dimensions typically
 comparable or some have larger than the
 cell thickness and depend on the initial
 field of 10–35 mT (Ref. 4 and 16) (Fig. 1)

the orientation of $\mathbf{M}(\mathbf{r})$ is opposite in-plane orientation, as depicted in Fig. 3(b), although $\mathbf{n}(\mathbf{r})$ and $\mathbf{M}(\mathbf{r})$ (Fig. 3(a)) because half-in-eigendirections are allowed in the $\mathbf{n}(\mathbf{r})$ line field but not in the $\mathbf{M}(\mathbf{r})$ vector field.

Although, in principle, simultaneous patterning of both $\mathbf{M}(\mathbf{r})$ and $\mathbf{n}(\mathbf{r})$ can be achieved by combining the two approaches described above, in practice, owing to the isolation of domain structure in $\mathbf{M}(\mathbf{r})$ when only the direction is patterned (Fig. 4). The $\mathbf{M}(\mathbf{r})$ within magnetic domains in the patterned region follows the spatial arrangement of the magnetic domains and behaves collectively (clipping opposite in-domain alignment). A applied magnetic field, the pol domain nature of the FLCC in explaining the topological required all connecting the half-in-eigendirection lines in $\mathbf{n}(\mathbf{r})$, causing a complete pattern of domains and all defects in explaining them, which solely depend on both the direction and strength of \mathbf{B} . In practice, the width of all defects within the region of direction or $\mathbf{n}(\mathbf{r})$ of even larger than half in region of uniform direction (Fig. 4). To counter the nature of magnetic in-domain alignment in the FLCC, the so-called dark field microscopy observation has revealed both location and orientation of individual nanoparticles (supplemental video S1)²² a zero field and when \mathbf{B} a different orientation, electrical charges, the domains of opposite \mathbf{M} . Unlike in conventional magnetic systems, these magnetic domains are typically

separated by the so-called Bloch or Néel walls⁸ in conventional localised collective deformations of $\mathbf{M}(\mathbf{r})$, magnetic in-domain alignment of the FLCC; indeed, no domain, so have the arrangement nature. This is because $\mathbf{M}(\mathbf{r})$ and $\mathbf{n}(\mathbf{r})$ are strongly coupled, so have collective deformation of $\mathbf{M}(\mathbf{r})$ between the domains would be coupled in terms of the corresponding elastic deformation of $\mathbf{n}(\mathbf{r})$. Indeed, the domain alignment in the FLCC has uniform direction but the $\mathbf{M}(\mathbf{r})$ so have here is no associated elastic free energy density, which all. A applied field, the in-domain alignment can be partially deprived of nanoparticles (supplemental video S1) and ranging in width from the arrangement being been individual nanoparticles of $\sim 1\mu\text{m}$, as determined by colloidal interaction between nanoparticles with different oriented dipole moments of the neighboring domains of opposite $\mathbf{M} \parallel \mathbf{n}_0$. When $\mathbf{n}(\mathbf{r}) \parallel \mathbf{M}(\mathbf{r})$ within the domains is directed, this in practice is further altered by the energetic cost of elastic deformation (Fig. 4). Although the domain alignment defects in the FLCC are arrangement in $\mathbf{M}(\mathbf{r})$, 2148/69 FDJ-0049/18, 4870(1)1327(1)

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